

## Short Communication

# NDDO MO Calculations

### IV. Correlation of *K*-Shell Binding Energy Shifts

Jeyaraman Chandrasekhar, Prem Kishore Mehrotra, Sankaran Subramanian, and Periakaruppan T. Manoharan

Structural Chemistry Group, Department of Chemistry, Indian Institute of Technology, Madras 600 036, India

Siegbahn's potential model as extended by Ellison *et al.* is used with density matrix elements calculated by the NDDO/2 procedure, to correlate the *K*-shell binding energy shifts of C, N and O atoms in a few molecules containing only the first-row atoms. The correlation is not superior to that obtained with the CNDO/2 method when only the monopole term is retained in calculating the Madelung potential energy. However, the results are in excellent agreement with experiment when the two-parameters model including the dipole and quadrupole terms is used.

**Key words:** *K*-shell binding energies

### 1. Introduction

X-ray photoelectron spectroscopy has become a powerful technique for studying molecular electronic structure and as a means of testing the quality of calculated molecular wave functions [1–3]. Several different approaches have been used to interpret the observed changes in core binding energies of particular atoms in different chemical environments, the so-called ESCA chemical shift. These vary in their sophistication, from simple correlation with the charge on the atom involved to direct calculation of the core binding energies [1–5]. A simple and popular model proposed by Siegbahn *et al.* [1] employs an expression of the type (1) for correlating the shifts.

$$E_B = kq + V + I. \quad (1)$$

In this equation  $q$  is the net charge on the atom under consideration,  $V$  is the potential due to electronic and nuclear charges on all other atoms.  $k$  and  $I$  are obtained by a least-squares procedure. This expression has been used in conjunction with CNDO/2 and INDO MO methods to predict the *K*-shell shifts for carbon,

nitrogen and oxygen atoms in various molecules [1], the average errors being around 1 eV for C and N shifts and 0.5 eV for O shifts.

Ellison *et al.* [6–7], while deriving the theoretical justification for this model, proposed two improvements: a two-parameter model in which the interaction of the core with the  $2s$  and  $2p$  orbitals on the same atom are treated individually, and inclusion of the dipole and quadrupole terms in calculating the extra-atomic Madelung potential term  $V$ . Within the CNDO/INDO framework the correlation improved considerably when the two-parameter model was used but the inclusion of dipole and quadrupole terms had little effect.

In the present series of papers [8–10] we have been considering in detail the potential of a modified version of Roby's NDDO MO [11, 12] method. Since the calculated wavefunctions provided a faithful representation of their *ab initio* counterparts, we have employed NDDO MO wavefunctions to correlate the  $K$ -shell binding energy shifts for Cl, N and O using the Siegbahn's model as extended by Ellison *et al.*

## 2. Method of Calculation

In the present model the shift in the  $K$ -shell binding energies as the electronic environment changes is written as

$$E_A = -k_s^A \Delta q_{As} - k_p^A \Delta q_{Ap} - \Delta V_A + l, \quad (2)$$

where  $\Delta q_{As} = q_{As} - q_{As}^0$  represents the differences between  $2s_A$  orbital electron population in the given molecule and in the reference molecule, and  $\Delta q_{Ap}$  is the corresponding difference between  $2p_A$  populations.  $k_s^A$  and  $k_p^A$  are parameters representing combinations of atomic Coulomb and exchange interaction between  $1s$  core and  $2s$  and  $2p$  orbital electrons on atom A respectively.  $l$  is a least-squares parameter to correct for errors inherent in the model.  $\Delta V_A = V_A - V_A^0$  is the difference in the Madelung potential energy at A arising from monopoles, dipoles, quadrupoles etc. centered on all atoms other than A, where

$$V_A = V_M + V_D + V_Q. \quad (3)$$

The monopole term:

$$V_M = \sum_{n \neq A} (q_n - Z_n^{\text{core}}) / R_{nA}. \quad (4)$$

The dipole term:

$$V_D = -5 \sum_{n \neq A} \sum_T P_{ST}^{(n)} T_{nA} / (3^{1/2} \zeta_n R_{nA}^3). \quad (5)$$

The quadrupole term:

$$V_Q = 3 \sum_{n \neq A} \sum_T \sum_U P_{TU}^{(n)} (3T_{nA} U_{nA} - R_{nA}^2 \delta_{TU}) / (2 \zeta_n^2 R_{nA}^5). \quad (6)$$

In the above equations,

$$q_n = \sum_r^n P_{rr} \quad (7)$$

represents the total charge on atom *n*. *T* and *U* are the Cartesian components *x*, *y*, *z*.  $T_{nA}$  is the *T*th coordinate of atom *n* relative to atom *A*.  $P_{sT}^{(n)}$  is the  $2s - 2p_T$  density matrix element.  $P_{TU}^{(n)}$  is the  $2p_T - 2p_U$  density matrix element.  $\zeta_n$  is the STO exponent on atom *n*.

Calculations were performed using both the single parameter model ( $k_s^A = k_p^A$  in Eq. 1) and the two-parameter model. In one set of calculations only the monopole term was included and in another  $V_D$  and  $V_Q$  terms were also included in evaluating  $V_A$ .

The density matrix elements over Löwdin symmetric orthogonalized orbitals were evaluated using the NDDO/2 scheme, the calculational details of which are described in [10]. While calculating the dipole and quadrupole terms it was assumed that the MO coefficients correspond to an STO basis rather than to the STO-3G basis actually employed in calculating the wavefunction. Also the exponents for C, N and O atoms were taken as 1.57, 1.92 and 2.23 respectively, values close to the Clementi-Raimondi *s* and *p* exponents [13] used in obtaining the STO-3G functions.

### 3. Results and Discussion

The  $1s$  binding energies were calculated using the various models for C, N and O atoms. The experimental values [1, 14] and the ones using a 2-parameter model including dipole and quadrupole terms are presented in Table 1. The average errors are also indicated in each case.

**Table 1.** Calculated  $1s$  ESCA shifts for C, N and O in some selected molecules using the two-parameters model with dipole and quadrupole terms

Carbon ( $1s$ ) <sup>a</sup>			Nitrogen ( $1s$ ) <sup>b</sup>			Oxygen ( $1s$ ) <sup>c</sup>		
Molecule	ESCA shift calculated	Exp. <sup>d</sup>	Molecule	ESCA shift calculated	Exp. <sup>d</sup>	Molecule	ESCA shift calculated	Exp. <sup>d</sup>
CH <sub>4</sub>	0.15	0.00	NH <sub>3</sub>	-0.11	0.00	H <sub>2</sub> O	-0.32	0.00
C <sub>2</sub> H <sub>4</sub>	0.28	-0.10	N <sub>2</sub>	4.26	4.35	CO <sub>2</sub>	1.76	1.44
C <sub>2</sub> H <sub>2</sub>	-0.15	0.40	CH <sub>3</sub> NH <sub>2</sub>	-0.33	-0.45	CH <sub>3</sub> OH	-0.53	-0.80
C <sub>2</sub> H <sub>6</sub>	0.05	-0.10	NNO	6.93	6.95	O <sub>2</sub>	3.57	3.84
CH <sub>3</sub> OH	1.84	1.90	NNO	2.97	3.05	N <sub>2</sub> O	1.53	1.54
CH <sub>2</sub> O	3.44	3.30	σ <sup>e</sup>	0.08		σ <sup>e</sup>	0.24	
CH <sub>3</sub> CHO	2.71	3.20						
CH <sub>3</sub> CHO	0.61	0.60						
CO <sub>2</sub>	7.10	6.84						
σ <sup>e</sup>	0.24							

<sup>a</sup>  $K_s = -6.984$ ,  $k_p = -15.564$ ,  $l = 0.151$

<sup>b</sup>  $K_s = -11.112$ ,  $k_p = -7.695$ ,  $l = -0.114$

<sup>c</sup>  $K_s = -21.206$ ,  $k_p = -21.301$ ,  $l = -0.317$

<sup>d</sup> = Ref. [1, 14]

<sup>e</sup> = Average error

The single parameter model with only the monopole term yields correlations with average errors around 1 eV in all three cases, similar to the results obtained with CNDO and INDO methods. On going over to the two-parameter model, the results, though not quite unsatisfactory, are still inferior to those obtained by the CNDO method at the same level of calculation. There has been a noticeable reduction in the average errors for O and N shifts but little in the case of C shifts.

The inclusion of the additional terms in  $V_A$ , however, has resulted in a steep fall in the average error. The corresponding two-parameter model has an average error of only 0.24 eV for C and O and 0.08 eV for N.

Although the number of molecules considered for N and O shifts is rather small, the excellent performance of NDDO/2 in the prediction of C shifts is indisputable. These results represent some of the best predictions of ESCA chemical shifts for C, N and O atoms by any MO method without the inclusion of valence relaxation effects. Taken together with the supposition of Ellison *et al.* [6] that the dipole and quadrupole terms in the extra-atomic Madelung potential will be important when calculated using better wavefunctions, it may be concluded that the NDDO/2 wavefunctions indeed provide faithful representations of *ab initio* functions.

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